

Electronic and Optical Excitations in Confined Nanostructures: Density Functional versus Many-body Methods

Serdar Ögüt

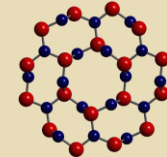
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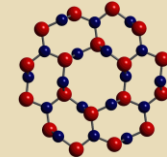


Supported by DOE under Grant No DE-FG02-03ER15488

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- **Introduction**
 - **Achievements and Deficiencies of DFT**
 - **Many-Body Theory Alternative to DFT**
- **GW & BSE for Electronic and Optical Excitations**
 - **Levels of Approximation**
 - **Excitons**
- **Examples (Benzene, Si_nH_m)**
- **Comparison of Many-Body vs DFT methods for Ag Clusters**
 - **ΔSCF vs GW**
 - **TDLDA vs BSE**



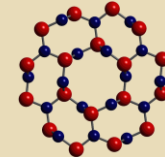
M. Tiago



J. C. Idrobo



K. Baishya



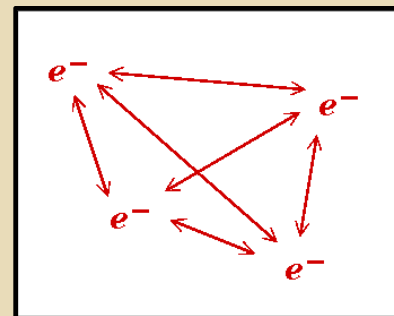
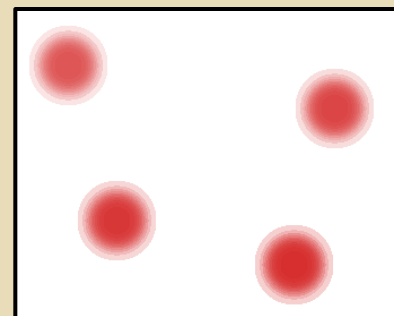
- Self-consistent mean-field theory.

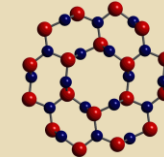
A generic electronic system is fully described by a single scalar function, the electron density $n(\mathbf{r})$.

- Formally exact.

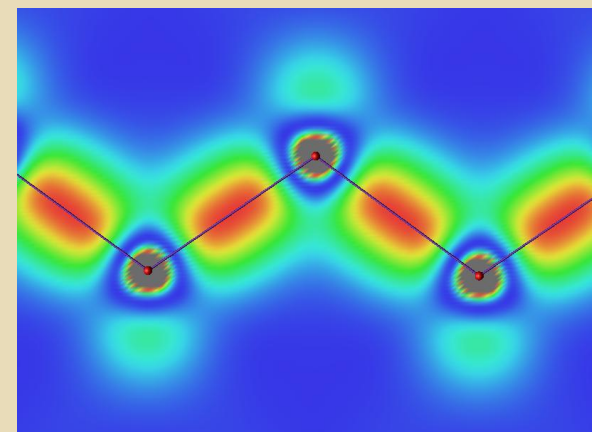
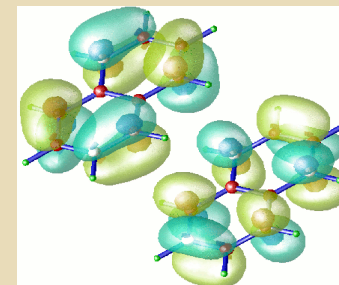
In practice, the exchange-correlation functional is modeled.

- Good description of weakly correlated, homogeneous systems.

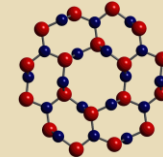
 $|\Psi\rangle$  $n(\mathbf{r})$



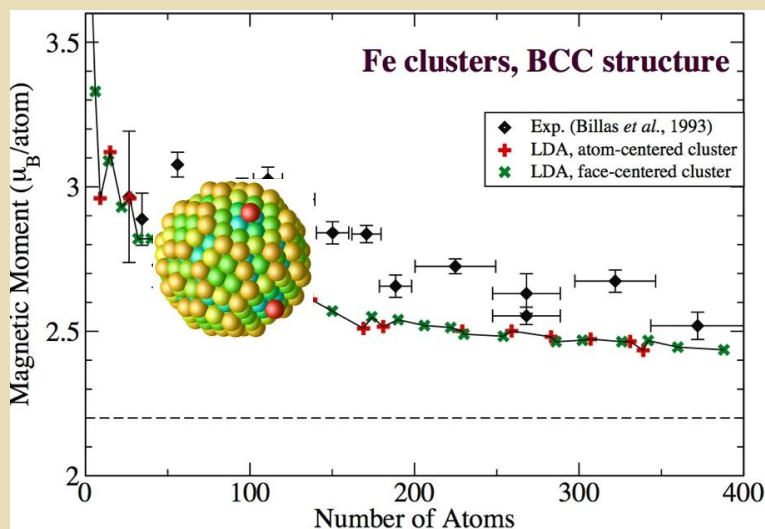
- First-principles electronic band theory.
- Quantum molecular dynamics.
- Quantum chemistry.



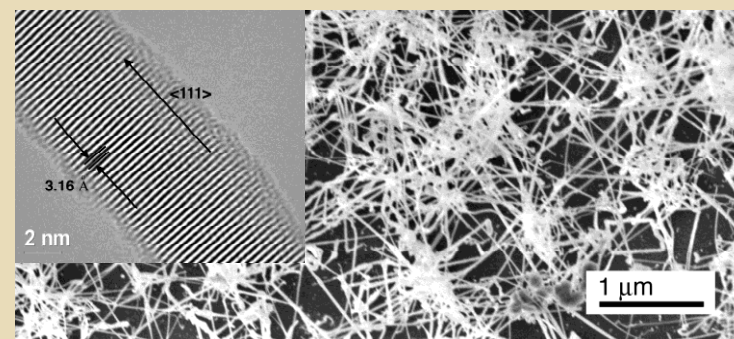
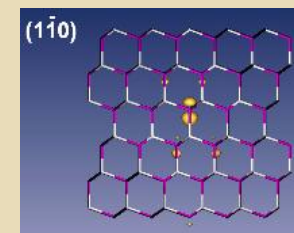
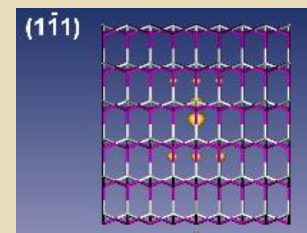
In 1998, Walter Kohn shared the Nobel Prize in Chemistry, “for his development of the density-functional theory”.



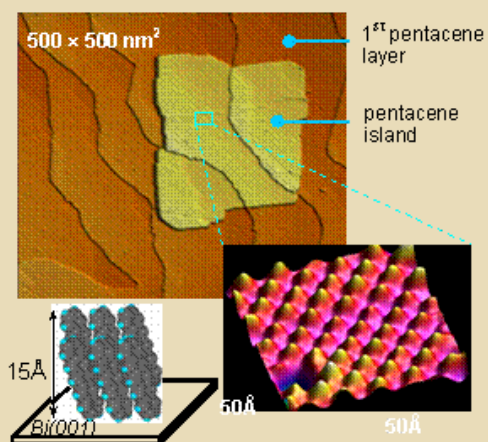
Magnetization in iron clusters



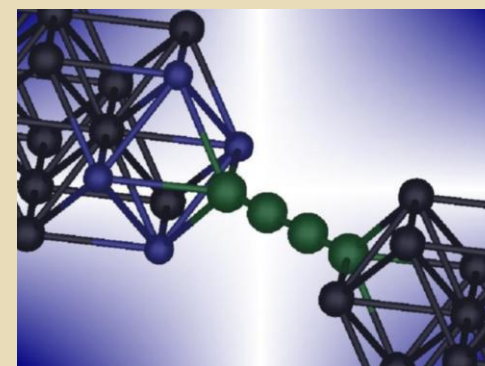
Doped semiconductor nanowires

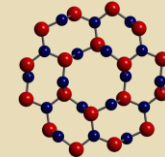


Growth of organic crystals

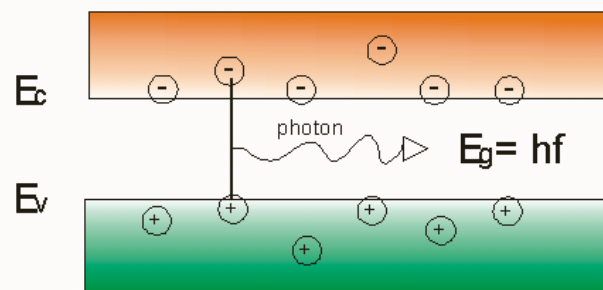


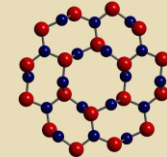
Transport through nanojunctions









- Exchange-correlation functional cannot be improved systematically.
- Complex magnetic ordering is still a challenge.
- Band gap is systematically underestimated.
- It does not give excitation energies: no optical spectroscopy.

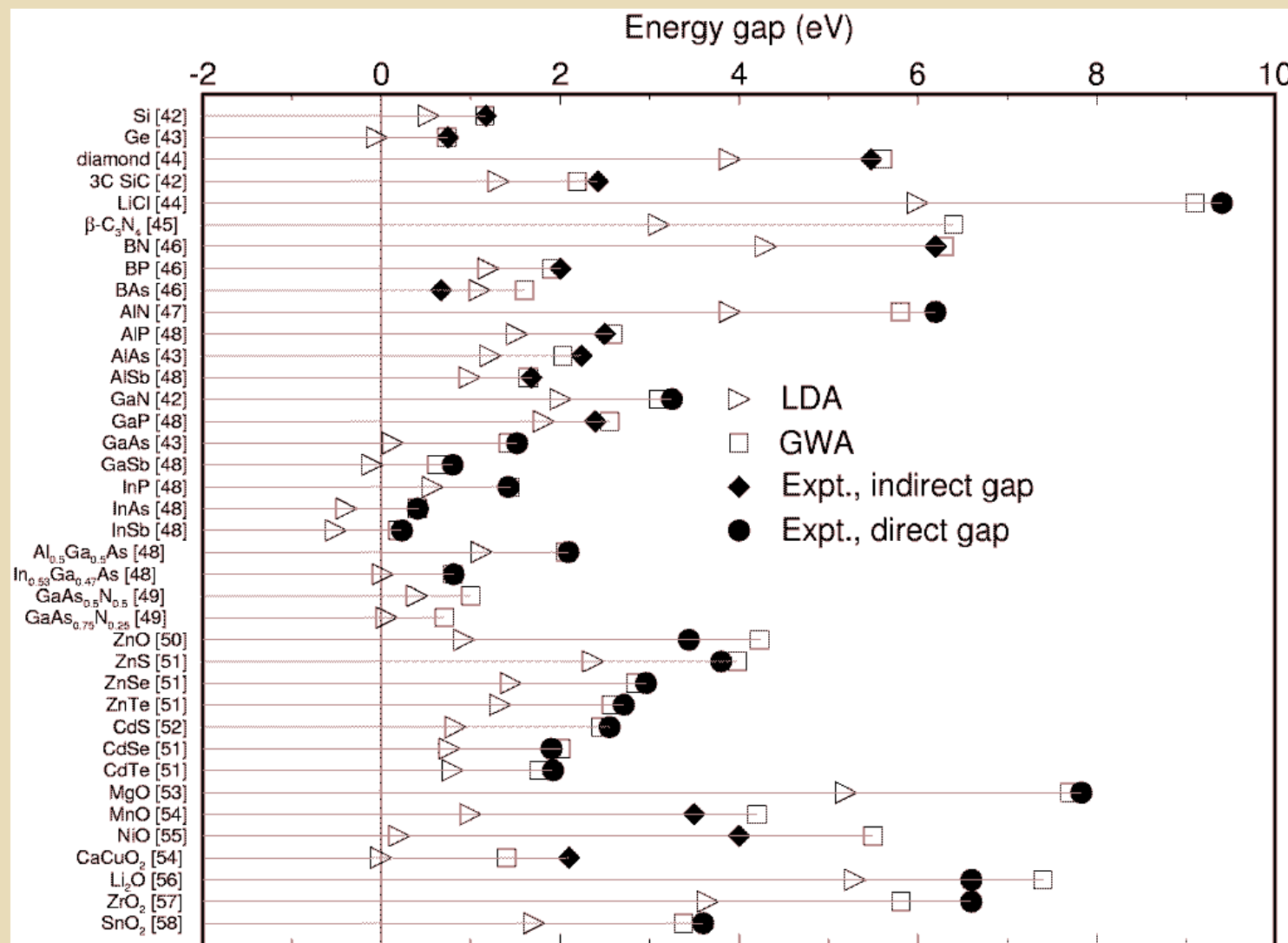
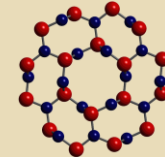




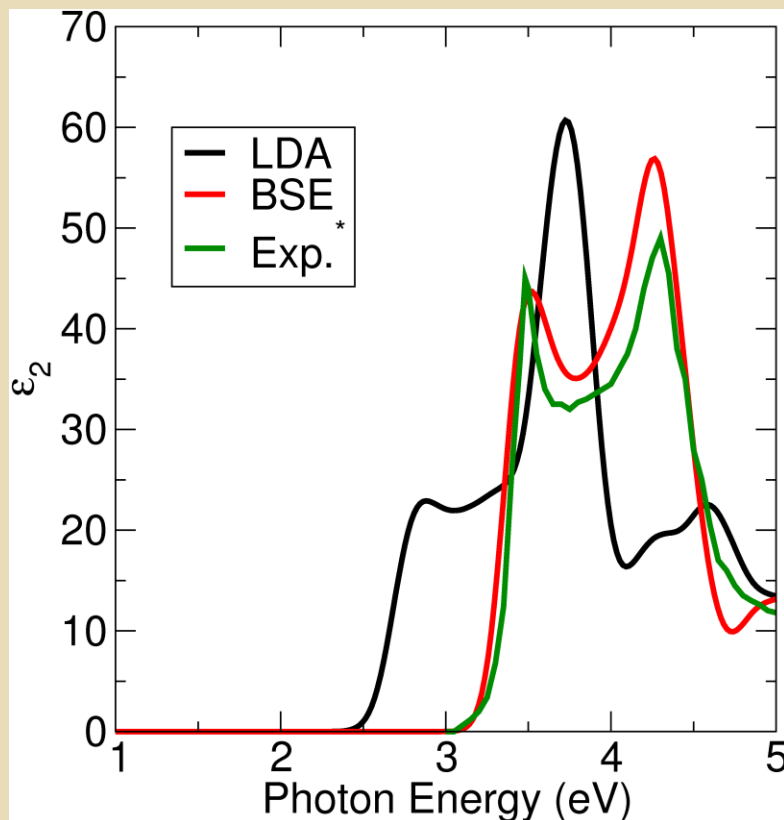
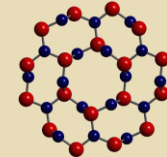
Yes! Many

Many-body Green's function methods

-  Gives band gap greatly improved from DFT.
-  Does not require phenomenological parameters.
-  Exact. If you manage to sum up all Feynman diagrams.
-  Amenable to high-performance computing.



Compiled by Aulbur, Jönsson, Wilkins, Solid State Phys. **54**, 1(2000)

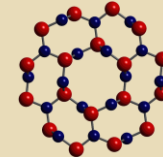


	E_{gap}	$\Gamma_{25'}^v \rightarrow \Gamma_{15}^c$	$\Gamma_{25'}^v \rightarrow X_1^c$
DFT	0.46	2.52	0.60
GW [†]	1.29	3.35	1.44
Exp.*	1.17	3.35	1.3

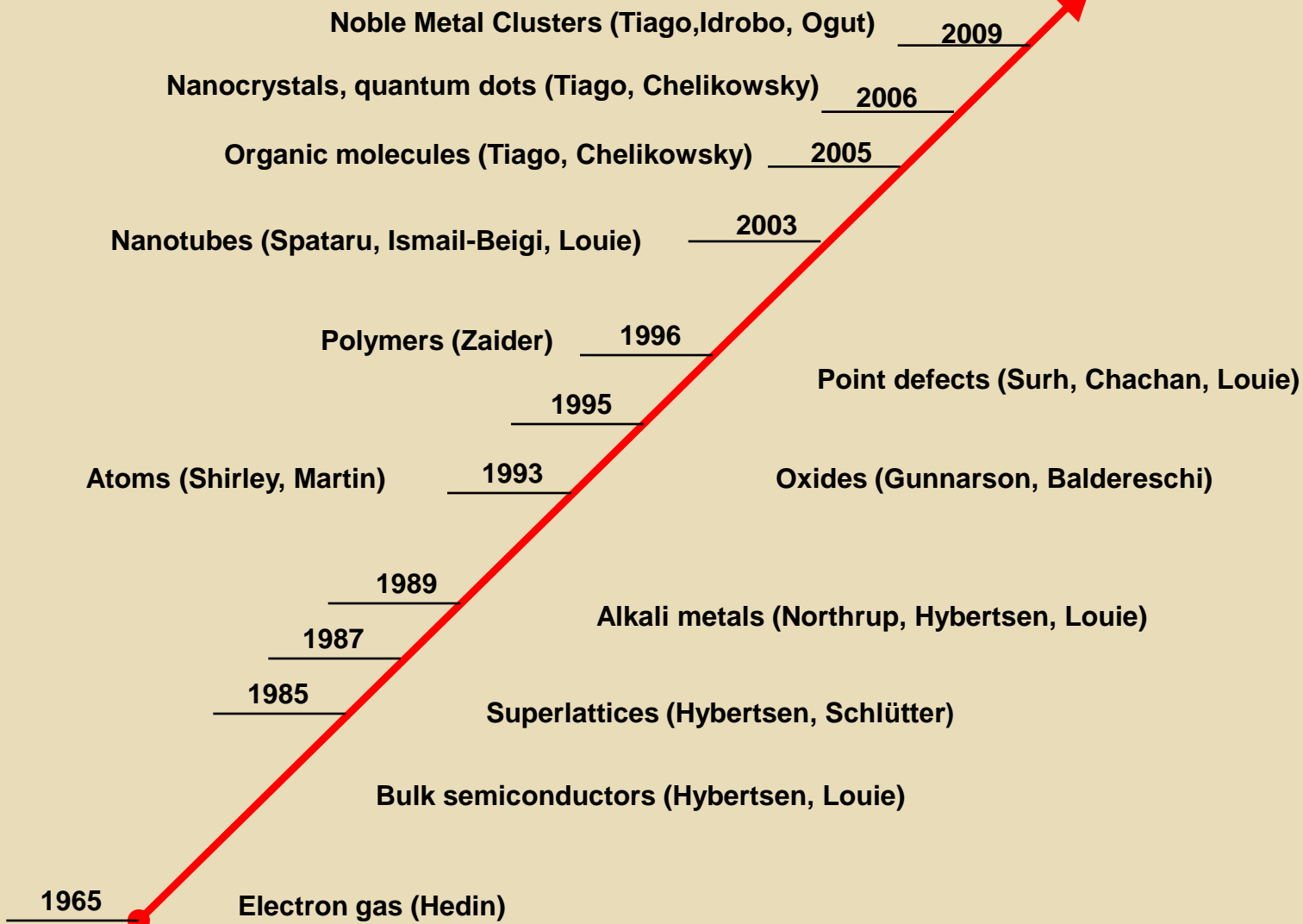
[†] Hybertsen and Louie (1986).

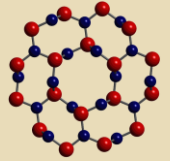
*Lautenschlager *et al.* (1987).

- Photoluminescence
- Exciton dynamics
- Photoemission spectroscopy
- Transport in doped materials



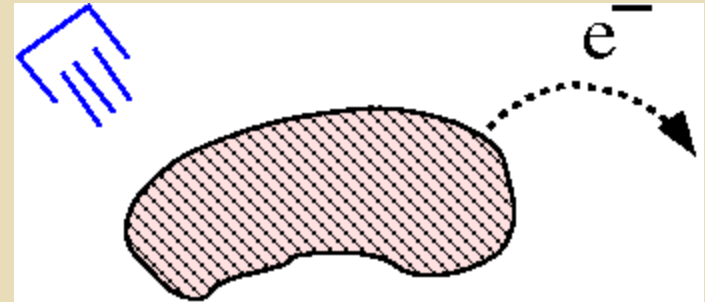
Biological systems, functionalized nanostructures, polarons, “dressed” excitons, photocatalysis ...





Quasi-particle approximation:

$$G(\mathbf{r}, \mathbf{r}'; E) = \sum_n \frac{\varphi_n(\mathbf{r}) \varphi_n^*(\mathbf{r}')}{E - \varepsilon_n}$$



Eigenvalue problem:

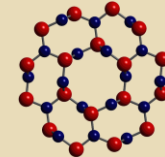
$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V_{ion} + V_H + \Sigma \right] \varphi_n(\mathbf{r}) = \varepsilon_n \varphi_n(\mathbf{r})$$

In DFT:

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V_{ion} + V_H + V_{xc} \right] \psi_n(\mathbf{r}) = E_n \psi_n(\mathbf{r})$$

Perturbation theory:

$$[H_{DFT} + \Sigma - V_{xc}] \varphi_n = \varepsilon_n \varphi_n$$



0-th order: “DFT”

$$\Sigma = V_{xc}$$

Hartree-Fock

$$\Sigma = iGV$$

1-st order: G_0W_0 approximation

$$\Sigma = iG_0W_0$$

1-st order + vertex: G_0W_f approximation

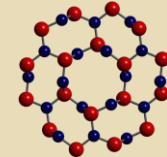
$$\Sigma = iG_0W_0\Gamma_0$$

Self-consistent: GW approximation

$$\Sigma = iGW$$

Exact: Hedin's equations

$$\begin{cases} W = V + VPW \\ P = -iGG\Gamma \\ \Sigma = iG\Gamma W \\ \Gamma = 1 + \frac{\delta\Sigma}{\delta G}GG\Gamma \end{cases}$$



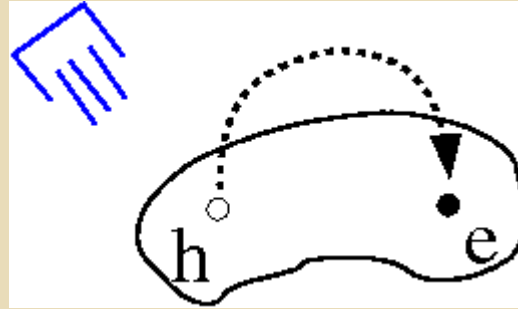
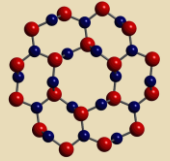
1. DFT Green's function:

$$G_0(\mathbf{r}, \mathbf{r}'; E) = \sum_n \frac{\psi_n(\mathbf{r})\psi_n^*(\mathbf{r}')}{E - E_n}$$

sum over occupied *and* unoccupied energy levels

2. G_0W_0 and G_0W_f approximations rely on DFT as a “good starting point”.
3. Local, dynamic dielectric screening must be known.

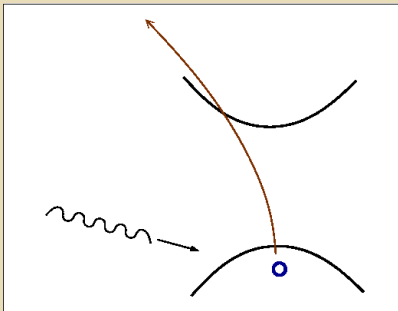
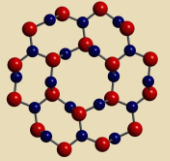
$$W \approx \epsilon^{-1} \frac{e^2}{|\mathbf{r} - \mathbf{r}'|}$$



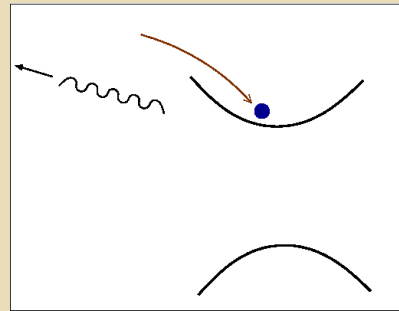
Eigenvalue problem:

$$\left(\epsilon_c^{GW} - \epsilon_v^{GW}\right) A_{cv}^S + \sum_{c'v'} \langle vc | K^{BSE} | c'v' \rangle A_{c'v'}^S = \Omega^S A_{cv}^S$$

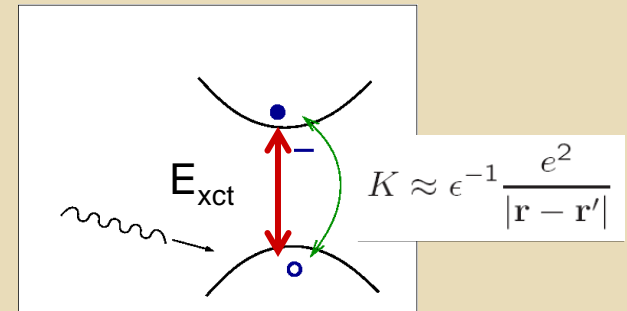
- Many-body expansion of the electron-hole propagator.
- Dynamics of electron-hole excitation obtained by solving the Bethe-Salpeter equation (BSE).
- Requires knowledge of quasiparticle orbitals (get from GW).



$$h\omega = -E_{\text{VBM}}$$

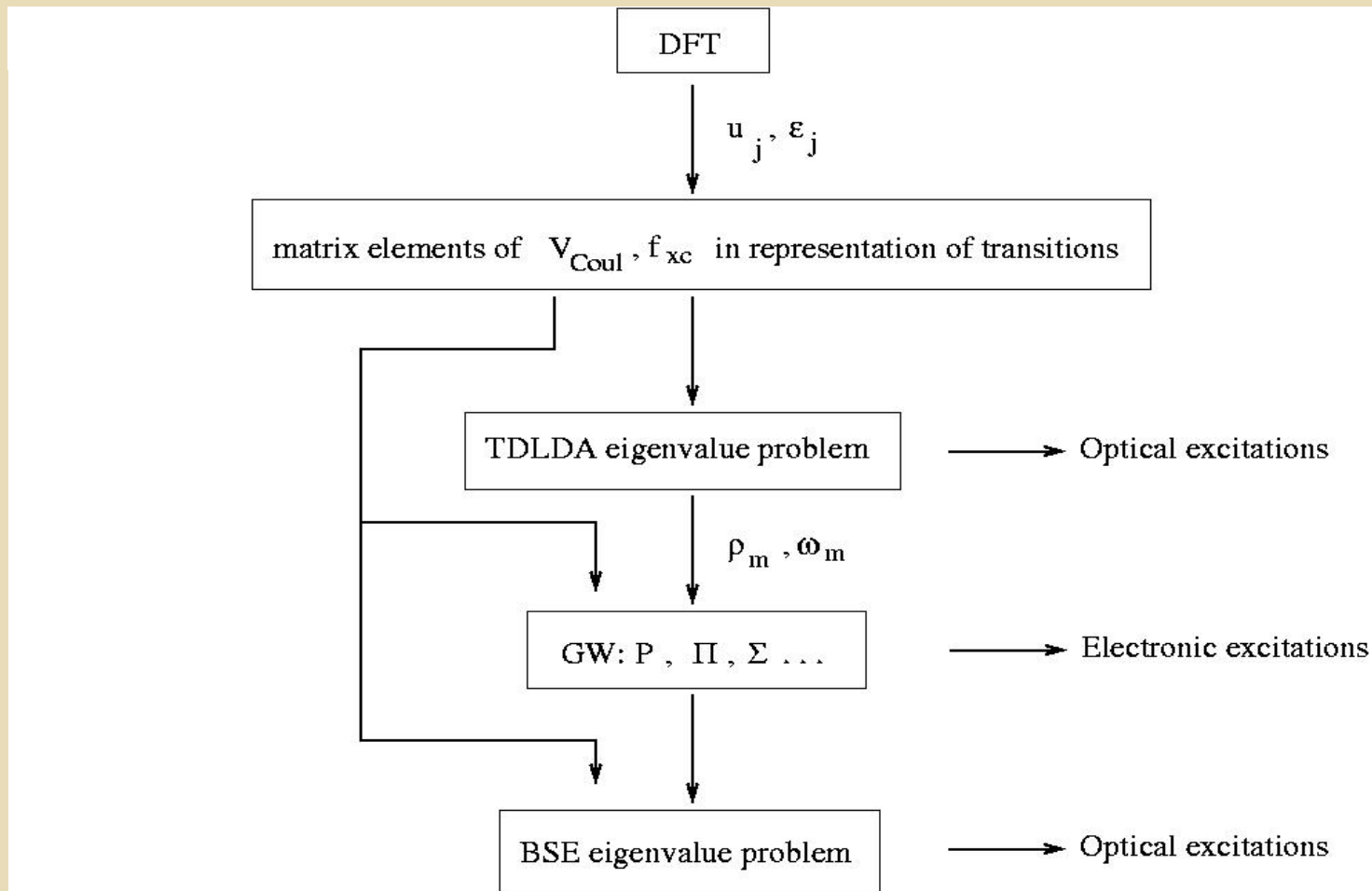
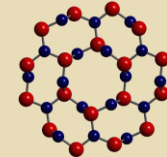


$$h\omega = -E_{\text{CBM}}$$

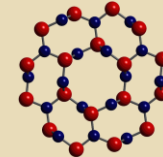
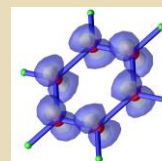


$$h\omega = E_{\text{xct}}$$

$$\text{Exciton binding energy : } (E_{\text{CBM}} - E_{\text{VBM}}) - E_{\text{xct}} > 0$$



RGWBS Package (free software)

1st Ionization Potential:

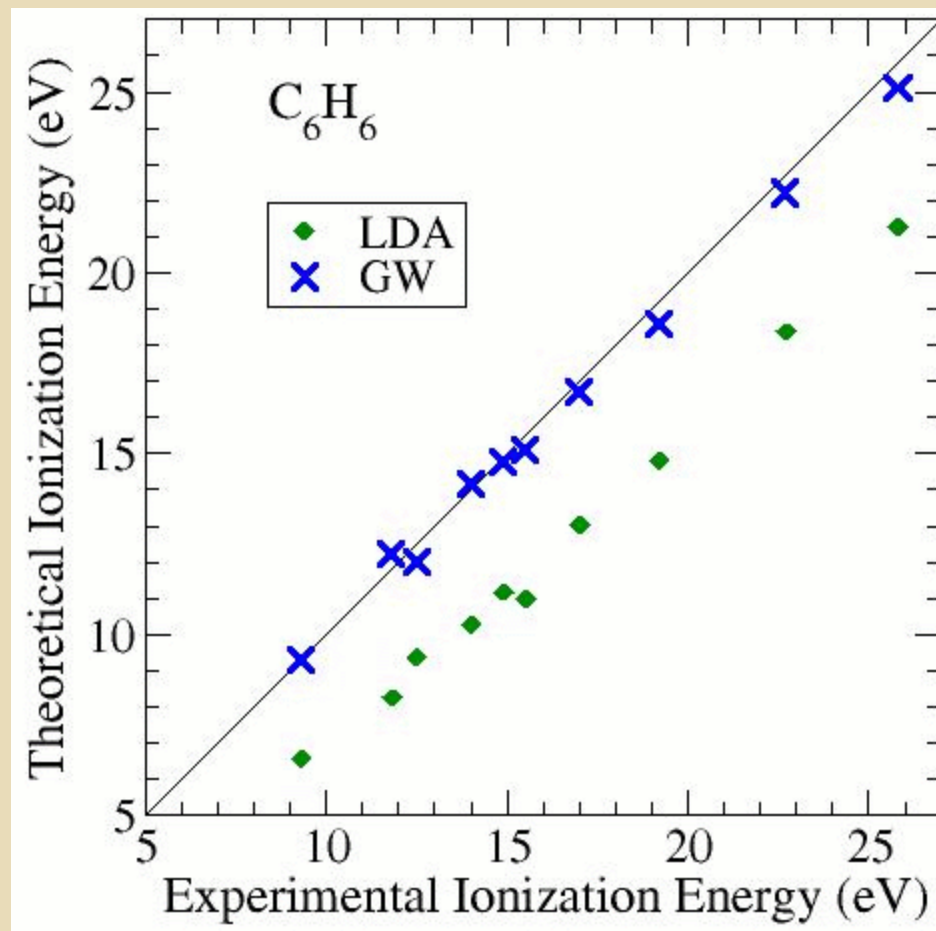
	IP (eV)
Δ SCF	9.55
GW	9.31
Exp.*	9.25

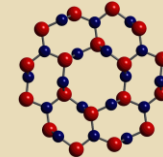
* Koch & Otto (1972).

Electron Affinity:

	EA (eV)
Δ SCF	-
GW	-0.99
Exp.*	-1.12

* Resonant transmission: Barrow *et al.* (1987).

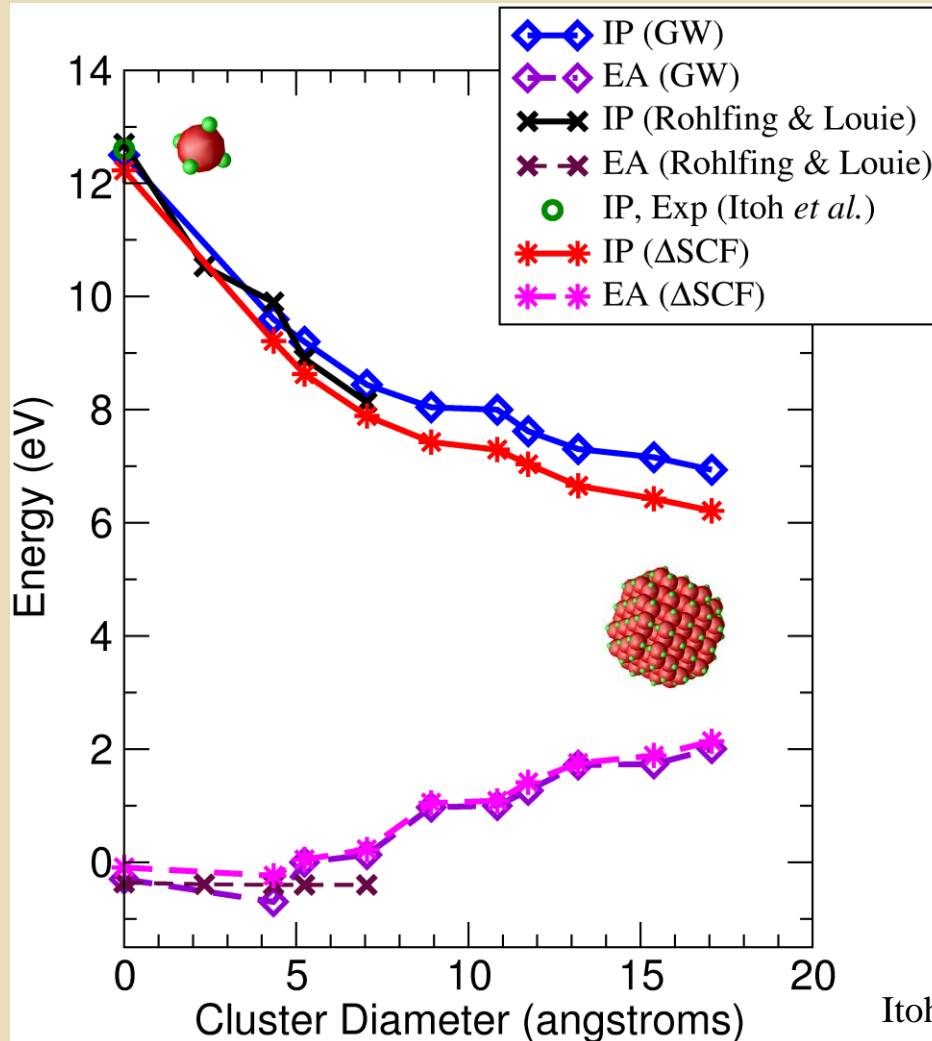
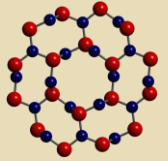




		TDLDA	BSE	Exp.* (eV)	
Triplet	B_{1u}	4.53	3.59	3.9	
Singlet	B_{2u}	5.40	4.86	5.0	dark
	B_{1u}	6.23	6.14	6.2	dark
	E_{1u}	6.9-7.2	7.23	6.9	bright

* Resonant transitions: Doering *et al.* (1969).

- Low-energy excitations are better described within BSE.
- Strong electron-hole binding is missing in TDLDA excitations.



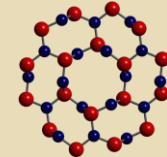
$$\Sigma = i \mathbf{G} \mathbf{W} \mathbf{T}_{\text{LDA}}$$

- Δ SCF predicts the correct EA.
- IP differs by ~ 0.7 eV:
 - self-interaction effects;
 - wrong asymptotic behaviour of V_{xc} .
- Same trend in Si nanoshells ($\text{Si}_{156}\text{H}_{184}$)

Rohlfing & Louie, Phys. Rev. Lett. **80**, 3320 (1998)

Itoh, Toyoshima & Onuki, J. Chem. Phys. **85**, 4867 (1986)

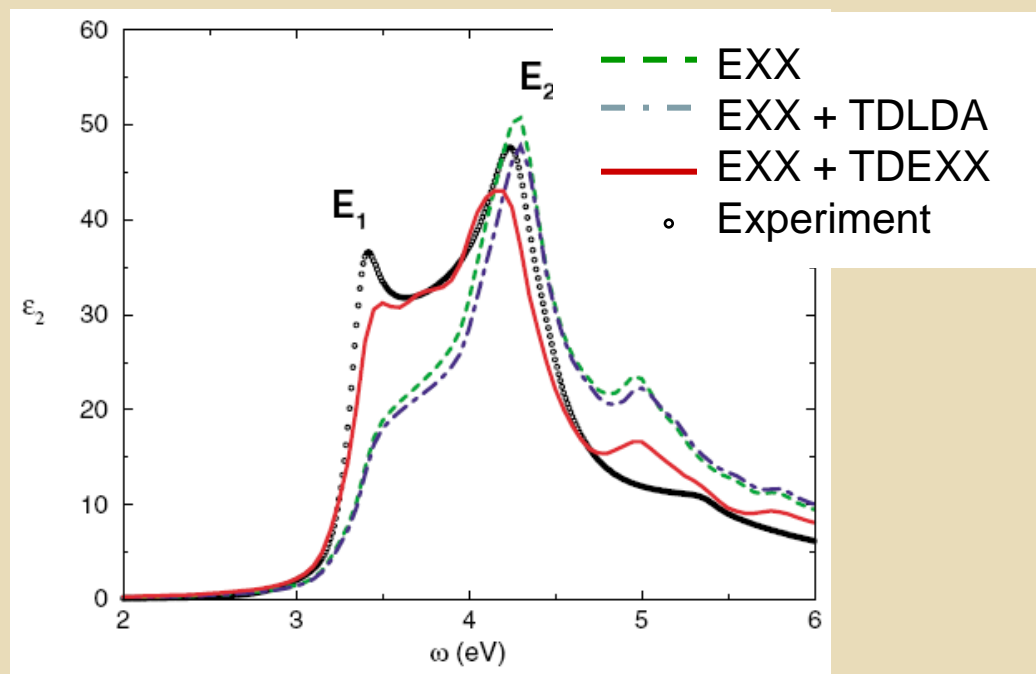
Tiago and Chelikowsky, Phys. Rev. B **73**, 205334 (2006)

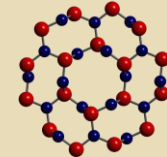


- Two of the state-of-the-art computational techniques for calculating optical excitations in materials: Time-dependent linear response theory using DFT (**TDDFT**) and adiabatic LDA (**TDLDA**), and Green's function many-body perturbation methods such as GW+Bethe-Salpeter Equation (**GWBSE**).

- Computational demand for TDLDA considerably smaller compared to GWBSE. GWBSE, however, gives much more accurate excitation energies in extended systems (excitonic effects).

Kim and Görling, PRL (2002)



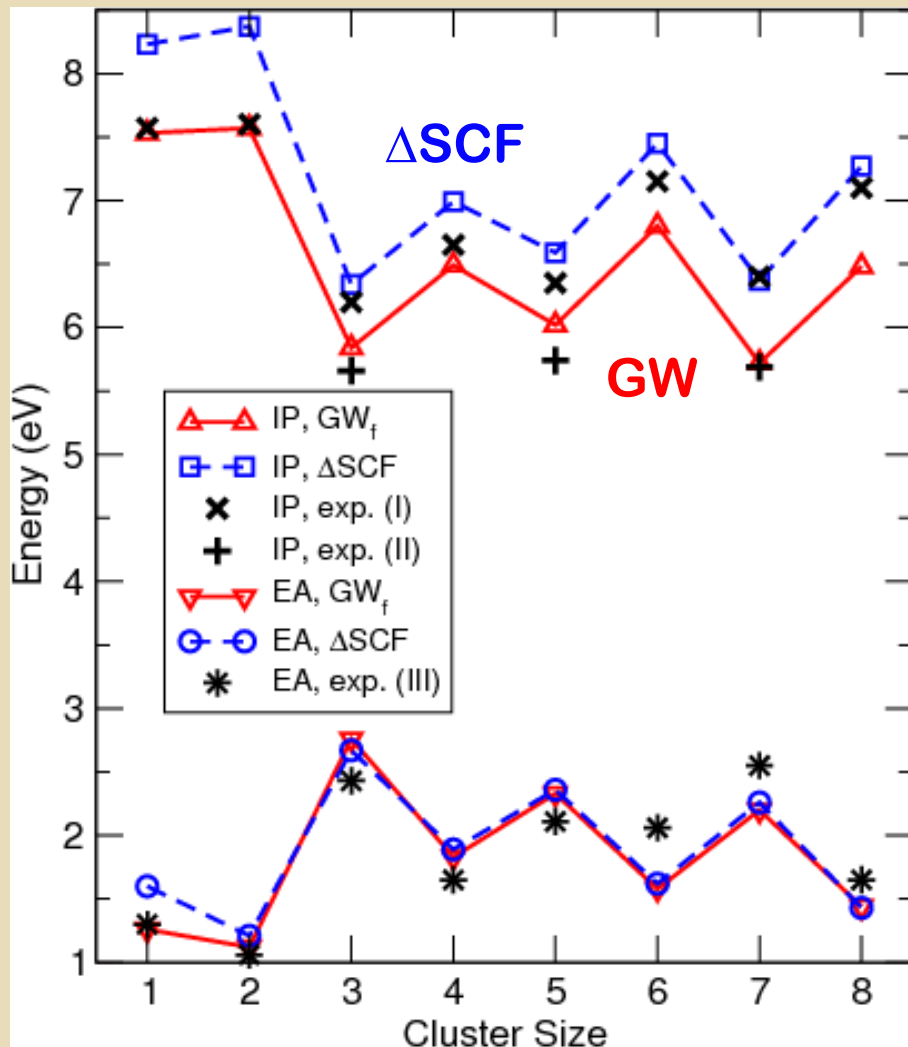
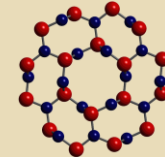


The two methods **mostly applied to *sp*-bonded clusters** with considerable success.

OUR GOAL

Systematic application and comparison of TDLDA and GWBSE in finite systems with **tightly bound *d* electrons to investigate the role of *d* electrons in optical excitations.**

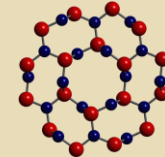
Chose Ag_n ($n = 1-8$) as prototype clusters with shallow (but filled) *d* orbitals. [Tiago, Idrobo, Ogut et al. PRB 79, 155419 (2009)]
Ongoing studies for Cu and Au clusters



- Within Δ SCF $IP = E(n-1) - E(n)$ and $EA = E(n) - E(n+1)$.
- Within GW $IP = -HOMO$ and $EA = -LUMO$.
- Generally quite **good agreement of GW results** with experiment (especially Ag and Ag₂). Agreement not so good with Δ SCF.
- Σ very sensitive to the number of virtual orbitals. Convergence accelerated by including a **static remainder** (estimate the numerical error by **truncating the sum over virtual orbitals at the level of COHSEX**).

	IP (eV)			EA (eV)		
	NS	ST	Exp	NS	ST	Exp
Ag	7.12	7.53	7.57	0.92	1.26	1.30
Ag ₂	6.27	7.54	7.60	0.82	1.12	1.06

(700 orbitals included in Σ)

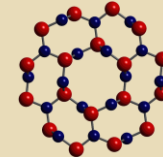


- Observed good agreement at this level of theory due to the fact that **HOMO and LUMO of Ag_n ($n \leq 8$) have almost entirely sp character** (little or no d character). For example, in Ag_2 , HOMO is **92% sp** , 8 % d , LUMO is **98 % sp** , and 2% d .
- Singly ionized Ag atom, Ag^+ , is a good test case for a system with **large (entirely) d character in HOMO (and purely s character in LUMO!)**. IP of Ag^+ is the double IP of Ag, which is experimentally available. EA of Ag^+ is the IP of neutral Ag!

	IP (eV)			EA (eV)		
	NS	ST	Exp	NS	ST	Exp
Ag^+	16.64	18.92	21.50	7.34	7.64	7.57

- **2.6 eV underestimate** of IP in Ag^+ ($4d$ level) due to **core-valence** separation in the pseudopotential construction (standard $4d^{10}5s^15p^0$ reference). Though $4s$ and $4p$ levels are ~ 80 and 50 eV below the $4d$ level, due to their **strong spatial overlap** with $4d$ levels, exchange and correlation among $4s$, $4p$, and $4d$ electrons are not described properly by a Slater type exchange-correlation [Rohlfing et al. PRL 75, (1995)]
- **Remedy: Create semi-core pseudopotentials** by keeping $4s$ and $4p$ levels in the valence. Use the reference $4s^24p^64d^{10}$ with sp core radii ~ 1.1 a.u. Very deep pseudopotentials. Use **$\hbar = 0.2$ a.u** (limited tests)

	NS	ST	Exp	NS	ST	Exp
$\text{Ag}^+(\text{semicore})$	20.67	21.85	21.50	7.04	7.30	7.57



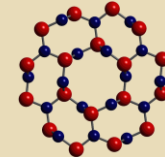
Ag

	TDLDA	GWBSE	Experiment
5s → 5p	4.09	3.68	3.74
5s → 6p	5.44	6.05	6.01

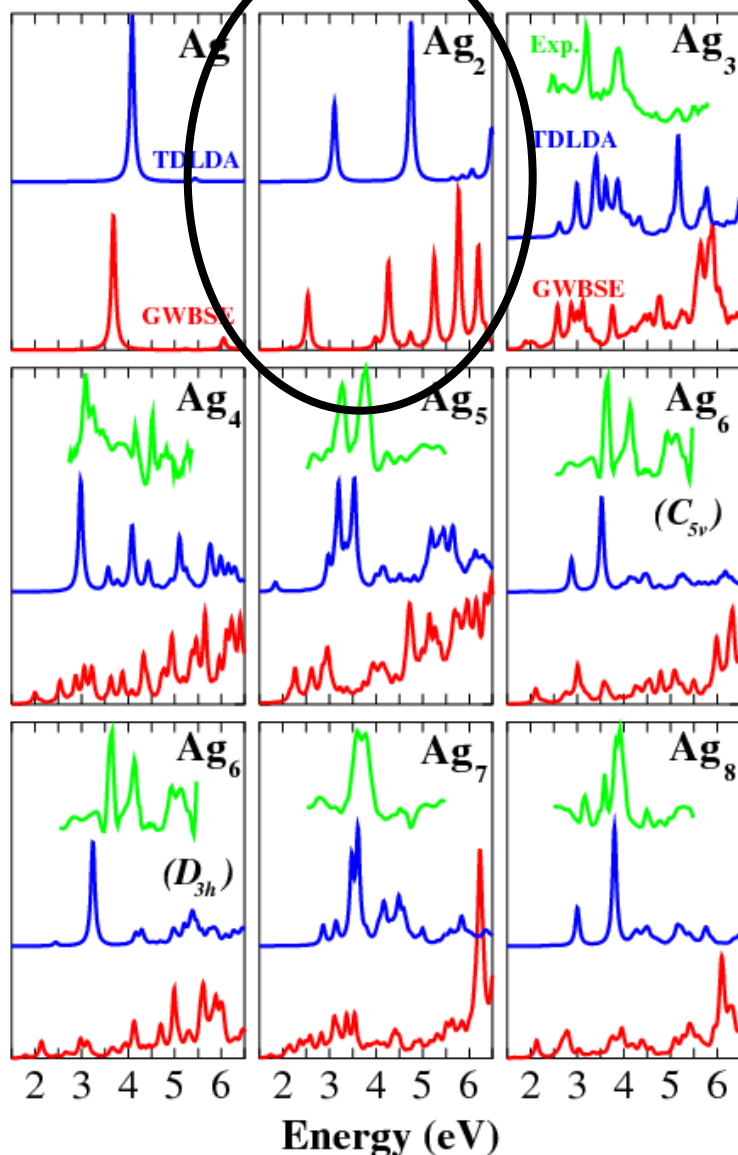
Ag₂

	TDLDA	GWBSE	Experiment
A - X	3.11	2.54	2.85
B - X	3.96	3.99	4.44
C - X	4.75	4.27	4.67

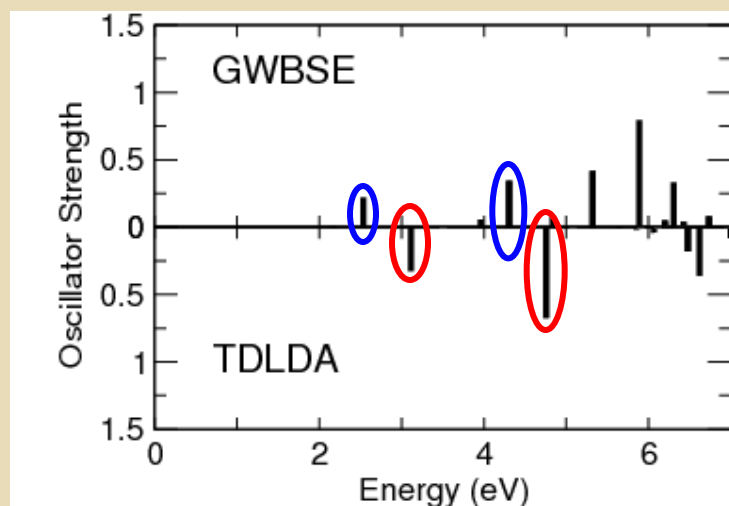
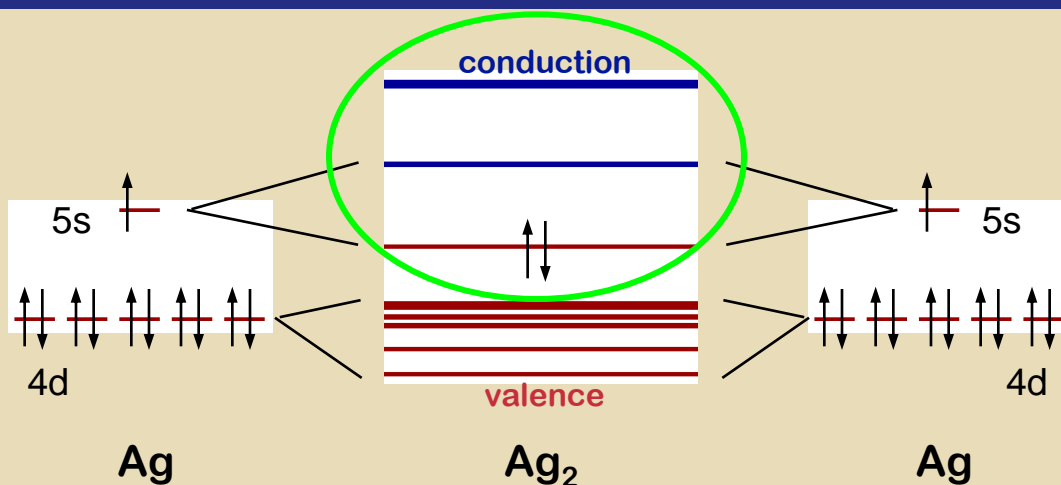
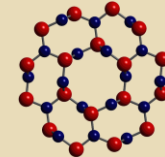
GWBSE clearly performs well for Ag atom. In Ag₂, agreement with experiment not as good. TDLDA is slightly better agreement (perhaps).



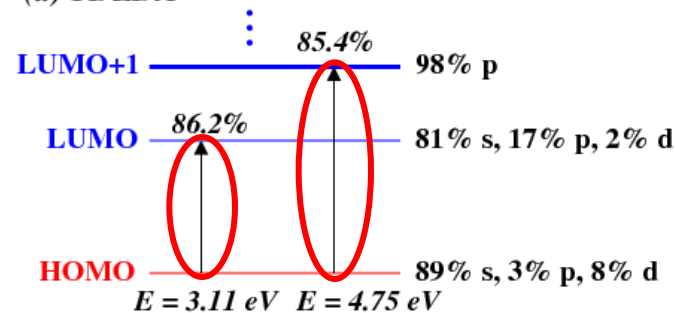
Absorption Cross Section (Arb. Units)



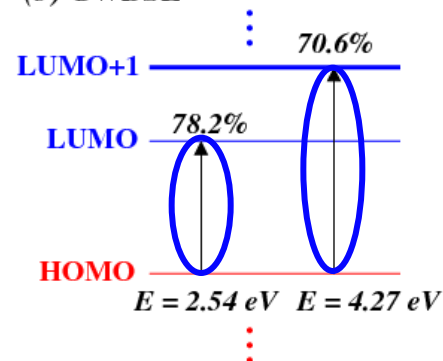
- For $n > 2$, **agreement between TDLDA and GWBSE very poor**. At low energies, OS from GWBSE quenched significantly, above 5 eV, high OS transitions.
- **TDLDA has clearly better agreement** with experimental data (esp. $n = 5 - 8$).
- Exchange-correlation effects involving $4d$ orbitals and strong non-locality of the BSE kernel are the main reasons for this behavior.
- Even a small mixture of d character results in **significant quenching** and **red-shifting** of the predicted transitions at the GWBSE level, as best illustrated for the case of Ag_2



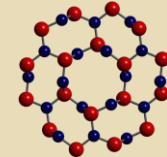
(a) TDLDA



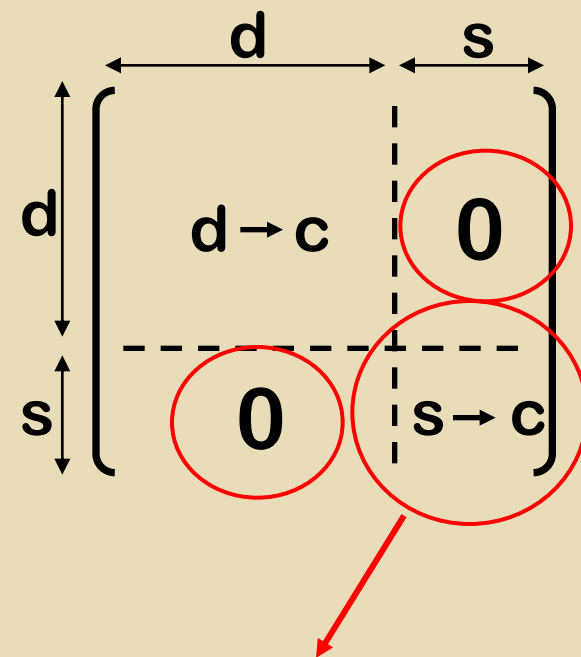
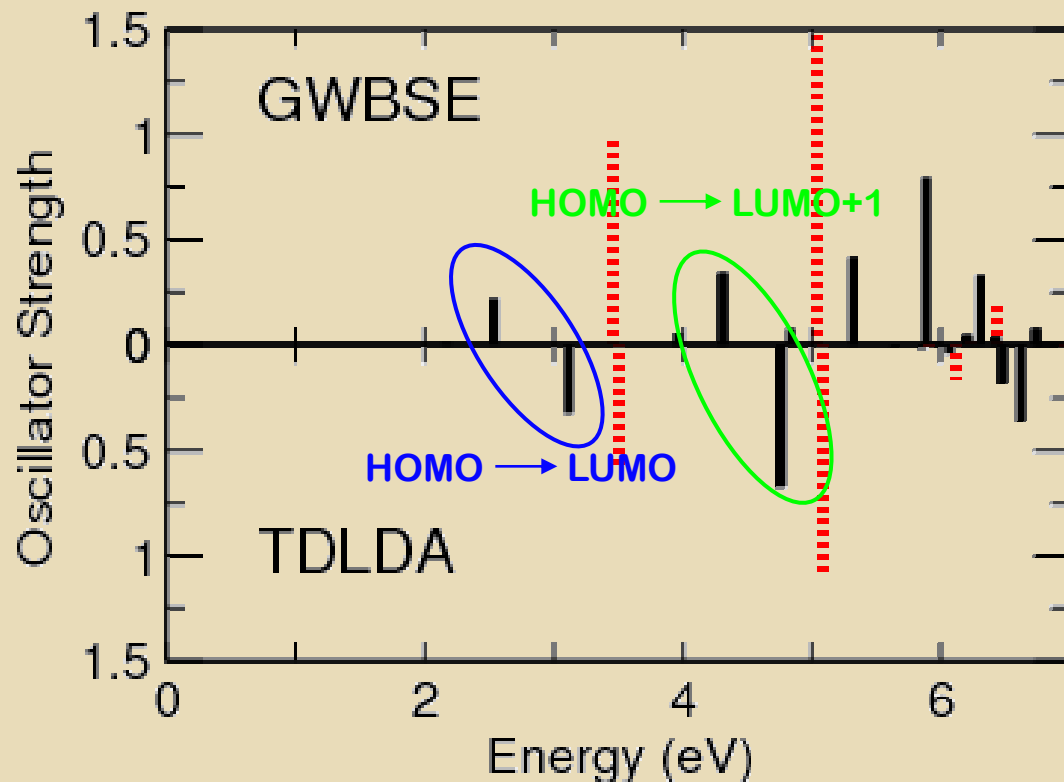
(b) GWBSE



The increase in d character of optical excitations at the GWBSE level accompanied by a redshift and quenching of OS compared to TDLDA.



$d: \nu = 1 - 10$ $s: \nu = 11$ (HOMO)



When d electrons are removed, very good agreement between TDLDA and GWBSE.